

Rotational and translational phonon modes in glasses composed of ellipsoidal particlesPeter J. Yunker,¹ Ke Chen,¹ Zexin Zhang,^{1,2,3} Wouter G. Ellenbroek,¹ Andrea J. Liu,¹ and A. G. Yodh¹¹*Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, Pennsylvania 19104, USA*²*Complex Assemblies of Soft Matter, Centre National de la Recherche Scientifique-Rhodia-University of Pennsylvania UMI 3254 Bristol, Pennsylvania 19007, USA*³*Center for Soft Condensed Matter Physics and Interdisciplinary Research, Soochow University, Suzhou 215006, China*

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The effects of particle shape on the vibrational properties of colloidal glasses are studied experimentally. “Ellipsoidal glasses” are created by stretching polystyrene spheres to different aspect ratios and then suspending the resulting ellipsoidal particles in water at a high packing fraction. By measuring displacement correlations between particles, we extract vibrational properties of the corresponding “shadow” ellipsoidal glass with the same geometric configuration and interactions as the “source” suspension but without damping. Low-frequency modes in glasses composed of ellipsoidal particles with major-to-minor axis aspect ratios of ~ 1.1 are observed to have predominantly rotational character. In contrast, low-frequency modes in glasses of ellipsoidal particles with larger aspect ratios (~ 3.0) exhibit a mixed rotational and translational character. All glass samples were characterized by a distribution of particles with different aspect ratios. Interestingly, even within the same sample it was found that small-aspect-ratio particles participate relatively more in rotational modes, while large-aspect-ratio particles tend to participate relatively more in translational modes.

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Although the “glass transition” occurs in a broad array of disordered systems, including molecular [1], polymer [2], granular [3], and colloidal glasses [4], much of the physics of granular and colloidal glasses has been derived from investigating ensembles of its simplest realization: spheres. The constituent particles of many relevant glasses, however, are anisotropic in shape or have orientation-dependent interactions. Such anisotropies are believed to affect many properties of glasses [5–10]. Thus, exploration of glasses composed of anisotropic particles holds potential for uncovering new consequences for both the physics of glasses and materials applications [11].

In glasses composed of frictionless spherical constituents, rotations of the spheres do not cost energy. Rotational modes therefore correspond to zero-frequency phonon excitations in the harmonic approximation. For anisotropic constituents, however, rotations are more energetically costly and can couple to translations. Glass vibrational properties, including the phonon density of states, are therefore expected to depend on the major-to-minor axis aspect ratio of constituent particles. Simulations of disordered systems with aspect ratios marginally greater than 1.0, for example, find low-energy rotational modes that are largely decoupled from translational modes [12,13]; apparently, when particles rotate in such systems, neighboring particles also rotate, but their positions remain essentially unperturbed.

Here we experimentally study glasses composed of ellipsoidal particles with aspect ratios α ranging from 1.0 to 3.0. To this end, we extend the displacement correlation matrix techniques employed in recent papers [14–17] to include rotations, and we employ video microscopy to derive the phonon density of states of corresponding “shadow” ellipsoidal glasses with the same geometric configuration and interactions as the experimental colloidal system but absent damping [15]. The spectra and character of the vibrational modes in these disordered media were observed to depend strongly on particle aspect ratio and particle aspect ratio *distribution*. For glasses

composed of particles with small median aspect ratios of ~ 1.1 , the lower-frequency modes are almost completely rotational in character, while higher-frequency ones are translational. In glasses of particles with larger aspect ratios (~ 3.0), significant mixing of rotations with translations is observed. In contrast to numerical findings for zero-temperature systems [12,13], we observe that the very lowest frequency modes for both systems have a mixed rotational-translational character, independent of aspect ratio. Additionally, even within the same sample, it was found that small-aspect-ratio particles tend to participate relatively more in rotational modes, while large-aspect-ratio particles tend to participate relatively more in translational modes. Evidently, the distribution of particle aspect ratios significantly affects phonon modes of glasses.

The experiments employ micron-sized polystyrene particles (Invitrogen) stretched to different aspect ratios [18–20]. Briefly, 3- μm -diameter polystyrene particles are suspended in a polyvinyl alcohol (PVA) gel and are then heated above the polystyrene melting point ($\sim 120^\circ\text{C}$) but below the PVA melting point ($\sim 180^\circ\text{C}$). In the process, the polystyrene melts, but the PVA gel only softens. The PVA gel is then placed in a vise and stretched. The spherical cavities that contain liquid polystyrene are stretched into ellipsoidal cavities. When the PVA gel cools, the polystyrene solidifies in the distorted cavities and becomes frozen into an ellipsoidal shape. The hardened gel dissolves in water, and the PVA is easily removed via centrifugation. Each iteration creates $\sim 10^9$ ellipsoidal particles in $\sim 50\ \mu\text{L}$ suspensions. Experiments are performed on samples stretched to 110% and 300% of their original size [snapshots of experimental particles are shown in the insets in Figs. 1(a) and 2(a)]. The stretching scheme produces a distribution of aspect ratios with a standard deviation of $\sim 18\%$. This distribution is most important for suspensions that are only slightly distorted from their initial spherical shape and therefore have greater propensity to crystallize. The distribution of aspect ratio $N(\alpha)$ for suspensions with more spherical particles [Fig. 1(a)] is peaked at $\alpha_{\text{peak}} = 1.1$, with

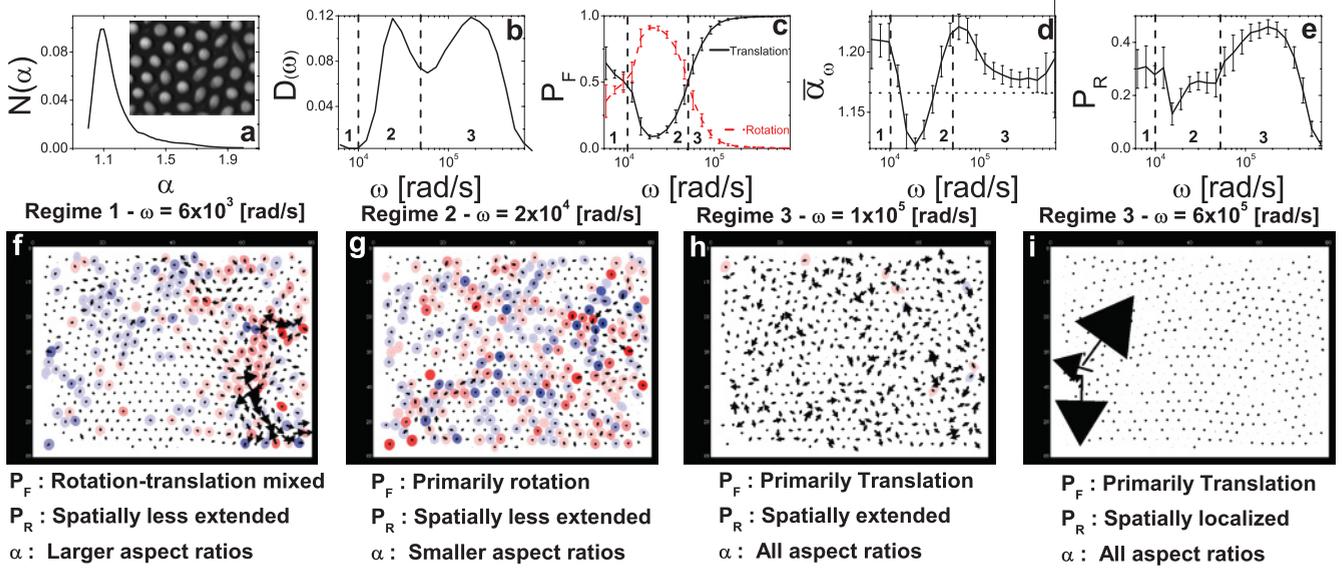


FIG. 1. (Color online) (a) Distribution of particle aspect ratio $N(\alpha)$ in samples with peak aspect ratio $\alpha_{\text{peak}} = 1.1$. The inset shows an experimental snapshot of part of the sample. (b) Vibrational density of states. Vertical dashed lines separate three distinct regimes corresponding to modes in the vector plots displayed in (f)–(i). (c) Translational (solid black line) and rotational (dashed red line) contributions to participation fraction P_F plotted versus frequency ω . (d) Participation-fraction-averaged aspect ratio $\bar{\alpha}_\omega$ plotted versus frequency ω . (e) Participation ratio P_R plotted versus frequency ω . (f)–(i) Displacement vector plots of eigenmodes from lowest frequency (f) to highest (i). The size of each arrow is proportional to the *translational* displacement of the particle at that position. The color (gray shading) intensity of each particle is proportional to the *rotational* displacement of the particle at that position (online, red indicates clockwise and blue indicates counterclockwise), with faint color (light gray) indicating small rotation and strong color (dark gray) indicating large rotation. Aspect ratio and frequency are specified in each plot.

mean aspect ratio $\bar{\alpha} = 1.2$, but it also has a long tail extending to $\alpha \sim 2.0$. A similar plot is shown in Fig. 2(a) for samples with $\alpha_{\text{peak}} = 3.0$ and $\bar{\alpha} = 3.3$.

Particles are confined between glass plates to quasi-two-dimensional chambers. From separate brightness calibration studies, we estimate the chambers to be no more than 5% larger than the minor axis particle length [15]. In all samples, dynamics are arrested (i.e., the average time it takes particles to move a distance greater than one tenth of the minor axis particle length is greater than our 10 000 s experimental window) (see Appendix A), and the spatial correlation functions of bond-orientational order decay exponentially (see Appendix B), with an average bond-orientational order parameter of 0.3 (0.03) for $\alpha_{\text{peak}} = 1.1$ (3.0). Nematic order is largely absent; the mean value of the nematic order parameter is 0.05, and the maximum value is 0.11 (see Appendix C).

Previous works have noted that the packing fraction at the jamming transition varies with particle shape [6]. In order to characterize how close our samples are to the jamming transition, we slowly evaporated water from the sample chamber. Complete evaporation should pack particles at the jamming transition for hard particles. We verified this claim for bidisperse mixtures of spheres with a size ratio of 1.4, where we find $\phi_{A,\text{MAX}} = 0.84(1)$, as expected. For ellipsoids with $\alpha_{\text{peak}} = 1.1$, $\phi_{A,\text{MAX}} = 0.87(1)$, consistent with [6,9,21]; the sample employed in this paper has $\phi_A = 0.86(1)$. For ellipsoids with $\alpha_{\text{peak}} = 3.0$, $\phi_{A,\text{MAX}} = 0.84(1)$, again consistent with [6,9,21]; the sample employed in this paper has $\phi_A = 0.83(1)$. Thus, both samples are near, but below, the jamming transition, with $\phi_{A,\text{MAX}} - \phi_A \approx 0.01$.

We extract vibrational properties by measuring displacement correlations. Specifically, we define $\mathbf{u}(t)$ as the $3N$ -component vector of the displacements of all particles from their average positions (\bar{x}, \bar{y}) and orientations $(\bar{\theta})[\mathbf{u}(t) = (x(t) - \bar{x}, y(t) - \bar{y}, \theta(t) - \bar{\theta})]$ and extract the time-averaged displacement correlation matrix (covariance matrix) $C_{ij} = \mathbf{u}_i \mathbf{u}_j^T$, where $i, j = 1, \dots, 3N_{\text{tot}}$ run over particles, positional and angular coordinates, and the average runs over time. In the harmonic approximation, the correlation matrix is directly related to the sample's stiffness matrix, defined as the matrix of second derivatives of the effective pair interaction potential with respect to particle position and angle displacements. In particular, $(C^{-1})_{ij} k_B T = K_{ij}$, where K_{ij} is the stiffness matrix. Experiments that measure C therefore permit us to construct and derive properties of a “shadow” ellipsoidal glass system that has the same static properties as our colloidal system (e.g., same correlation matrix and same stiffness matrix, but no damping) [15]. Following [22], we expect undamped hard particles that repel entropically, near but below the jamming transition, to give rise to solidlike vibrational behavior on time scales that are long compared to the collision time but short compared to the time between particle rearrangement events [14,17]. The stiffness matrix arising from entropic repulsions is directly related to the dynamical matrix characterizing vibrations $D_{ij} = \frac{K_{ij}}{m_{ij}}$, where $m_{ij} = \sqrt{m_i m_j}$ and m_i is an appropriate measure of inertia. For translational degrees of freedom, $m_i = m$, where m is the particle mass. For rotational degrees of freedom, $m_i = I_i$ represents the particle moment of inertia with respect to axes centered about each particle's center of mass and pointing in

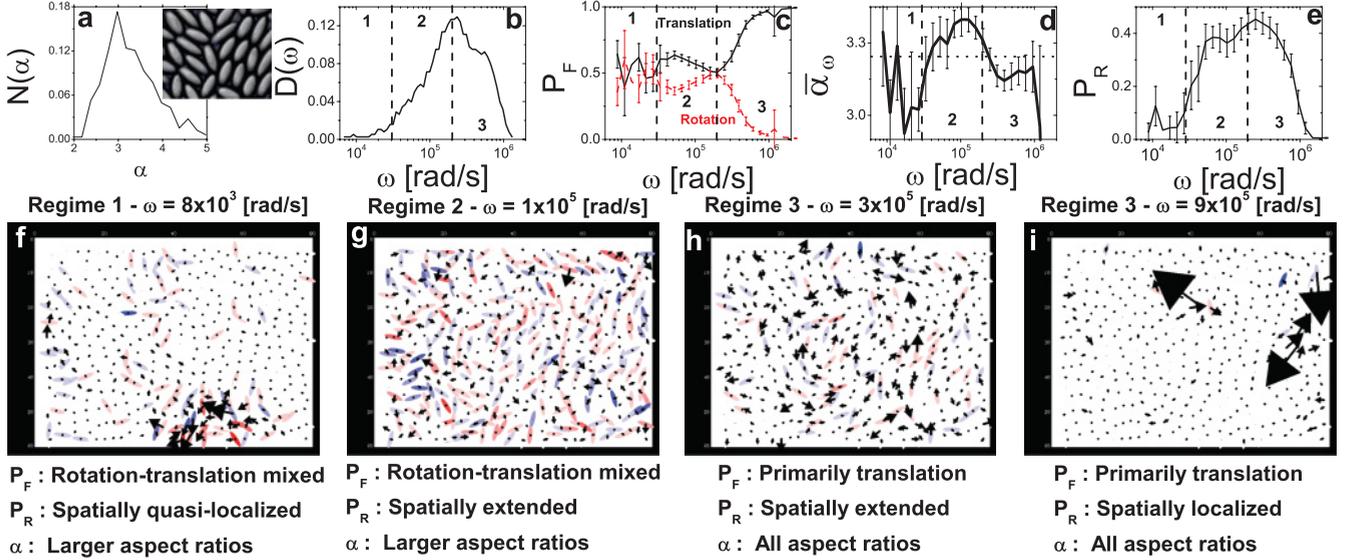


FIG. 2. (Color online) (a) Distribution of aspect ratio $N(\alpha)$ in samples with peak aspect ratio $\alpha_{\text{peak}} = 3.0$. The inset shows an experimental snapshot of part of the sample. (b) Vibrational density of states. Vertical dashed lines separate three distinct regimes corresponding to modes in the vector plots displayed in (f)–(i). (c) Translational (solid black line) and rotational (dashed red line) contributions to participation fraction P_F plotted versus frequency ω . (d) Participation-fraction-averaged aspect ratio $\bar{\alpha}_\omega$ plotted versus frequency ω . (e) Participation ratio P_R plotted versus frequency ω . (f)–(i) Displacement vector plots of eigenmodes from lowest frequency (f) to highest (i). The size of each arrow is proportional to the *translational* displacement of the particle at that position. The color (gray shading) intensity of each particle is proportional to the *rotational* displacement of the particle at that position (online, red indicates clockwise and blue indicates counterclockwise), with faint color (light gray) indicating small rotation and strong color (dark gray) indicating large rotation. Aspect ratio and frequency are specified in each plot.

the z direction; $I_i = m(a_i^2 + b_i^2)/2$, where a_i and b_i are the major and minor radii of the i th ellipsoid. The eigenvectors of the dynamical matrix correspond to amplitudes associated with the various phonon modes, and the eigenvalues correspond to the frequencies/energies of the corresponding modes. Data were collected over 10 000 s so that the number of degrees of freedom, $3N \approx 2000$, is small compared to the number of time frames (~ 8000) [15]. Additionally, we find K_{ij} is far above the noise only for adjacent particles, as expected.

The vibrational density of states $D(\omega)$ is plotted in Fig. 1(b) for the system with $\alpha_{\text{peak}} = 1.1$. $D(\omega)$ exhibits two distinct peaks. Zero-temperature simulations find that these peaks split completely for α sufficiently close to 1 and for sufficiently small systems close enough to the jamming transition [12,13]. For ellipsoids with $\alpha_{\text{peak}} = 3.0$ [Fig. 2(b)], on the other hand, $D(\omega)$ has a single peak, consistent with numerical predictions [12,13]. Thus, the vibrational spectrum of ellipsoids with small anisotropy is significantly different from those of spheres and those of ellipsoids with higher aspect ratio.

To quantitatively explore the modes, we calculated several different quantities. We will first introduce all of these quantities and then discuss them all at the same time. First, to quantify the translational and rotational contributions to each mode, we sum the participation fractions P_F of translational and rotational vibrations over all particles for each mode. The eigenvectors of each mode are normalized such that $\sum_{m,n} e_\omega(m,n)^2 = 1$, where m runs over all particles and n runs over all coordinates. The participation fraction for particle m , component n , in mode with frequency ω is then $P_F(\omega) = e_\omega(m,n)^2$. Thus, the *translational* participation fraction in a mode with frequency ω is

$P_{F,XY}(\omega) = \sum_{m=1,\dots,N,n=X,Y} e_\omega(m,n)^2$, and the *rotational* participation fraction is $P_{F,\theta}(\omega) = 1 - P_{F,XY}(\omega) = \sum_{m=1,\dots,N} e_\omega(m,\theta)^2$. Translational and rotational participation fractions are plotted in Figs. 1(c) and 2(c).

To investigate the effects of aspect ratio polydispersity, we measure the eigenvector-weighted ellipsoid aspect ratio as a function of mode frequency. Specifically, we compute $\bar{\alpha}_\omega = \sum_{m,n} \alpha_m e_\omega(m,n)^2$, where α_m is the measured aspect ratio of particle m . Thus, $\bar{\alpha}_\omega$ is a measure of the average particle aspect ratio for the particles participating in mode ω [Figs. 1(d) and 2(d)].

Finally, to assess the degree of mode localization, we quantify the spatial extent of individual modes by computing the participation ratio $P_R(\omega) = [\sum_{m,n} e_\omega(m,n)^2]^2 / [N_{\text{tot}} \sum_{m,n} e_\omega(m,n)^4]$ [Figs. 1(e) and 2(e)]. The participation ratio provides an indication of mode localization in space. If a mode is localized, a small number of terms will dominate, making $\sum_{m,n} e_\omega(m,n)^4$ and $[\sum_{m,n} e_\omega(m,n)^2]^2$ similar in size, so $P_R(\omega) \approx 1/N$.

Representative modes are shown in Figs. 1(f)–1(i) and 2(f)–2(i) for samples with $\alpha_{\text{peak}} = 1.1$ and 3.0, respectively. Modes from all samples fall qualitatively into three regimes. For $\alpha_{\text{peak}} = 1.1$, three distinct regimes exhibiting different behavior are labeled in Figs. 1(b)–1(e). For mode frequencies higher than $\omega \approx 54\,000$ rad/s, i.e., frequencies above the “dip” separating the two peaks in the density of states [Fig. 1(b)], the modes (regime 3) are translational in character. Interestingly, the lowest-frequency modes in regime 3 are spatially extended [Fig. 1(h)], while the highest-frequency modes are spatially localized [Fig. 1(i)], similar to the modes in glasses composed of spheres. Modes just above $\omega \approx 54\,000$ rad/s are enriched

in longer ellipsoids and have a mixed translational-rotational character. Modes in regime 2, extending from $1300 \lesssim \omega \lesssim 54\,000$ rad/s, are strongly rotational in character and are concentrated on small-aspect-ratio particles [Fig. 1(g)]. In regime 1, below $\omega \approx 1300$ rad/s, modes again have a mixed rotational-translational character and are concentrated on longer particles [Fig. 1(f)]. Regime 1 was not observed in numerical simulations of monodisperse ellipsoid packings at zero temperature [12,13]. We also find that the mean value of elements of the stiffness matrix connecting particles to their neighbors decreases as aspect ratio increases (see Appendix D); this observation suggests that longer ellipsoids are more weakly coupled to their neighbors and are relatively more likely to be excited at low frequency.

Figures 2(b)–2(e) show that for $\alpha_{\text{Peak}} = 3.0$, high-frequency modes above $\omega \approx 3 \times 10^5$ rad/s in regime 3 are translational in character, with a nearly average mode-averaged aspect ratio, resembling those of spheres. These translational modes cross from extended [Fig. 2(h)] to localized [Fig. 2(i)] at the upper end of the spectrum. Modes with $20\,000 \lesssim \omega \lesssim 3 \times 10^5$ rad/s in regime 2 are extended, with a mixed rotational-translational character, and are slightly enriched with longer ellipsoids at higher frequencies and shorter ellipsoids at somewhat lower frequencies [Fig. 2(g)]. In regime 1, $\omega \lesssim 2 \times 10^4$ rad/s, modes are again slightly enriched in larger-aspect-ratio particles and are quasilocalized with mixed translational-rotational character [Fig. 2(f)].

Comparing the two systems, the behaviors of modes at high frequencies (regime 3) and at the lowest frequencies (regime 1) are qualitatively very similar. The largest qualitative differences between large- and small-aspect-ratio systems occur in regime 2, where modes have primarily rotational character for systems with $\alpha_{\text{Peak}} = 1.1$ and modes have mixed translational-rotational character for systems with $\alpha_{\text{Peak}} = 3.0$.

To summarize, experiments suggest that the nature of low-frequency modes in glasses depends strongly on constituent particle aspect ratio. Rotational modes tend to occur at lower frequencies than translational vibrations, and for glasses with aspect ratios of ~ 1.1 , a frequency regime exists wherein the spectrum is strongly rotational in character, consistent with numerical results [12,13]. Additionally, even within each sample, particles with small aspect ratios tend to participate more in rotational modes, while particles with larger ones tend to participate more in translational modes. We also find low-frequency modes enhanced in larger aspect ratio particles with mixed rotational-translational character that were not present in simulations. The *distribution* of particle aspect ratio $N(\alpha)$ is thus an important physical factor affecting phonon modes of ellipsoidal glasses. Recent work suggests that low-participation-ratio, low-frequency modes appear to correlate with regions prone to rearrangement or plastic deformation [23]. Thus, the existence of additional low-frequency modes concentrated around particles with certain aspect ratios may have important consequences for the mechanical response of glasses.

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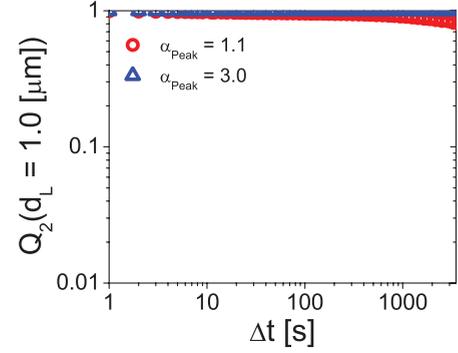


FIG. 3. (Color online) The two-point correlation function Q_2 , which probes self-overlap on the preselected length scale d_L , is plotted versus delay time for ellipsoidal glasses with different aspect ratios. Dynamic arrest is apparent.

APPENDIX A: GLASSY DYNAMICS

As a first step toward elucidation of glass dynamics in these systems, we compute the two-time self-overlap correlation function: $Q_2(d_L, \Delta t) = \frac{1}{N_{\text{tot}}} \sum_{i=1}^{N_{\text{tot}}} \exp\left(-\frac{\Delta r_i(\Delta t)^2}{2d_L^2}\right)$ (Fig. 3) [24]. Here d_L is a preselected length scale to be probed, N_{tot} is the total number of particles, and $\Delta r_i(\Delta t)$ is the distance particle i moves in time Δt . If a particle moves a distance smaller than d_L , Q_2 will remain approximately unity; if a particle moves a distance greater than d_L , Q_2 will fall to zero. Notice that for glasses of each aspect ratio, $Q_2(d_L = 1.0 \mu\text{m})$ decays very little over the experimental time scale, thereby indicating that glass dynamics are arrested at length scales of order of the particle size.

APPENDIX B: BOND-ORIENTATIONAL ORDER

To demonstrate the absence of long-range orientational order in these systems, the bond-orientational order parameter $\psi_6 = \frac{1}{N_{\text{tot}} z} \sum_{j=1}^{N_{\text{tot}}} \left| \sum_{k=1}^z e^{i6\theta_{jk}} \right|$ and its spatial correlation function $g_6(r) = \langle \psi_6^*(r_i) \psi_6(r_j) \rangle$ are calculated (Fig. 4). Here θ_{jk} is the angle between the x axis and the $j-k$ bond between particles j and k , z is the coordination number of particle j , and r_i and r_j are the positions of particles i and j . g_6 decays faster in samples with $\alpha_{\text{Peak}} = 3.0$ than it does in

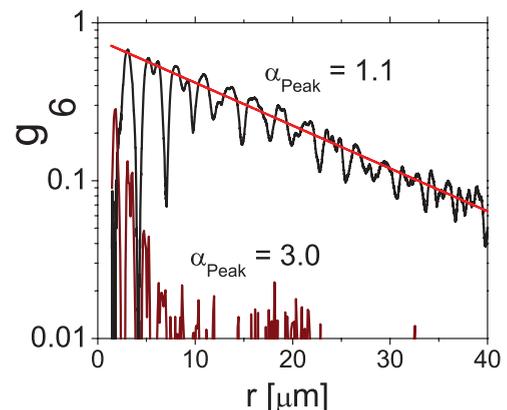


FIG. 4. (Color online) Bond-orientational order spatial correlation functions $g_6(r)$ for ellipsoidal glasses with different aspect ratios.

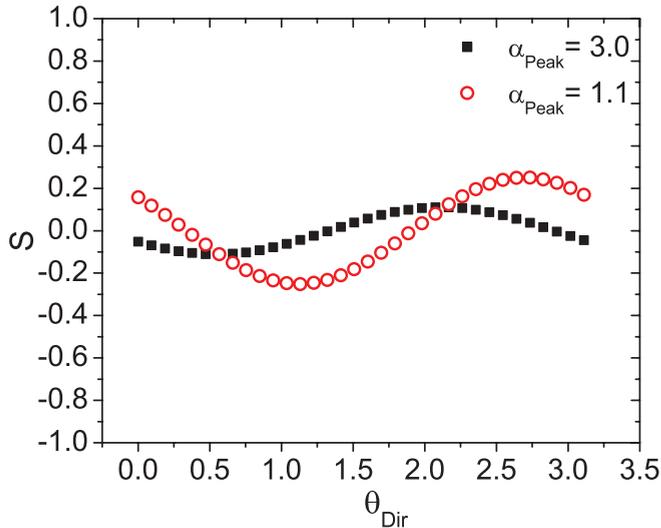


FIG. 5. (Color online) The average nematic order parameter S as a function of the director angle for samples with $\alpha_{\text{Peak}} = 3.0$ (solid squares) and $\alpha_{\text{Peak}} = 1.1$ (open circles).

samples with $\alpha_{\text{Peak}} = 1.1$. However, g_6 decays exponentially in each sample (see exponential fit line in Fig. 4), a signature of structural disorder characteristic of glasses (e.g., [25]).

APPENDIX C: NEMATIC ORDER

To demonstrate the absence of long-range nematic order in these systems, the nematic order parameter, $S = \sum_{j=1}^{N_{\text{tot}}} 2 * \cos(\theta_j - \theta_{\text{Dir}})^2 - 1$, where θ_j is the orientation of particle j and θ_{Dir} is the orientation of the nematic director, and angle brackets represent ensemble averaging, is calculated (Fig. 5). For an isotropic distribution of orientations, $S = 0$, and for perfectly aligned particles, $S = 1$. The mean value of S in our large-aspect-ratio samples ($\alpha_{\text{Peak}} = 3.0$) is 0.05, and the maximum value of S is 0.11. The mean value of S in samples with $\alpha_{\text{Peak}} = 1.1$ is 0.00, and the maximum value of S is 0.25.

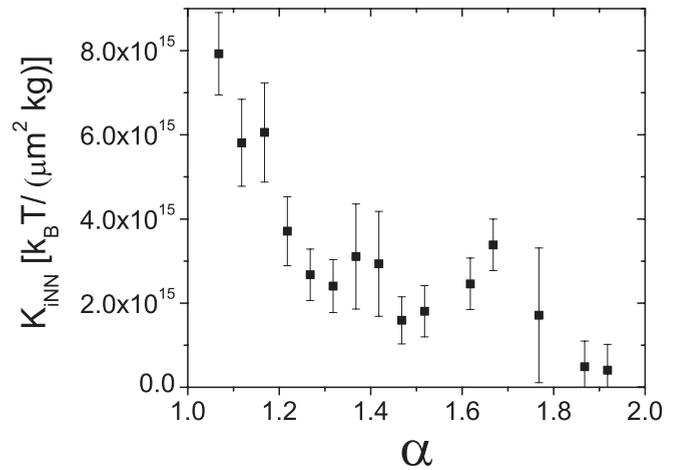


FIG. 6. The average spring constant K_{iNN} connecting nearest neighbors as a function of aspect ratio α for samples with $\alpha_{\text{Peak}} = 1.1$. Error bars represent standard error.

APPENDIX D: LOW-FREQUENCY MODES WITH MIXED ROTATIONAL-TRANSLATIONAL CHARACTER

Low-frequency modes for samples with $\alpha_{\text{Peak}} = 1.1$ have mixed rotational-translational character. These modes were not seen in zero-temperature simulations in which all particles have identical aspect ratios [12,13]. These “mixed” modes typically involve larger-aspect-ratio particles. To understand why these modes appear at low frequencies, we calculated the average spring constant connecting a particle’s rotation to its nearest neighbors $K_{iNN} = \langle K_{ij} / m_{ij} \rangle_{NN}$, where $\langle \rangle_{NN}$ indicates an average over nearest-neighbors pairings, i runs over all θ components, and j runs over all components. We then plotted K_{iNN} as a function of aspect ratio (Fig. 6). K_{iNN} decreases as α increases, indicating that the average spring constraining rotation decreases as α increases. Smaller spring constants K_{iNN} lead to vibrations at smaller frequencies. Thus, particles with larger aspect ratios tend to vibrate at lower frequencies.

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